

# Fourier Inversion of the Star Transform

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Radiology/Bioengineering/Applied  
Math & Computational Science

Joint work with

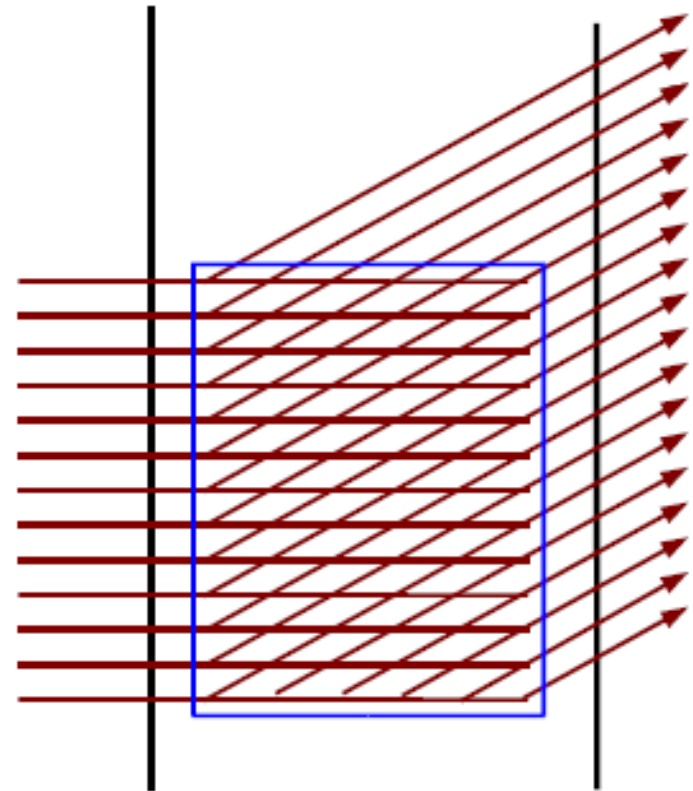
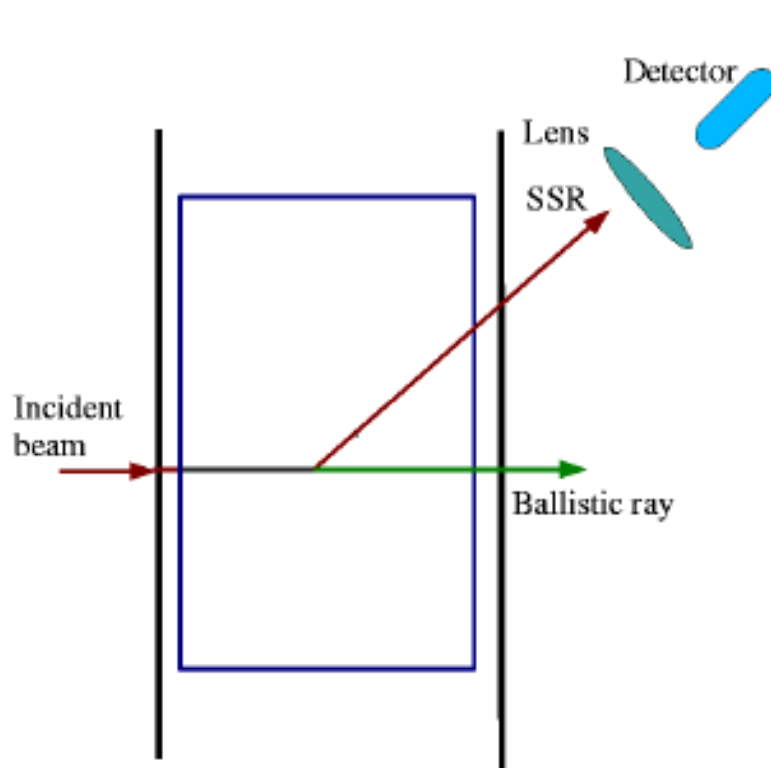
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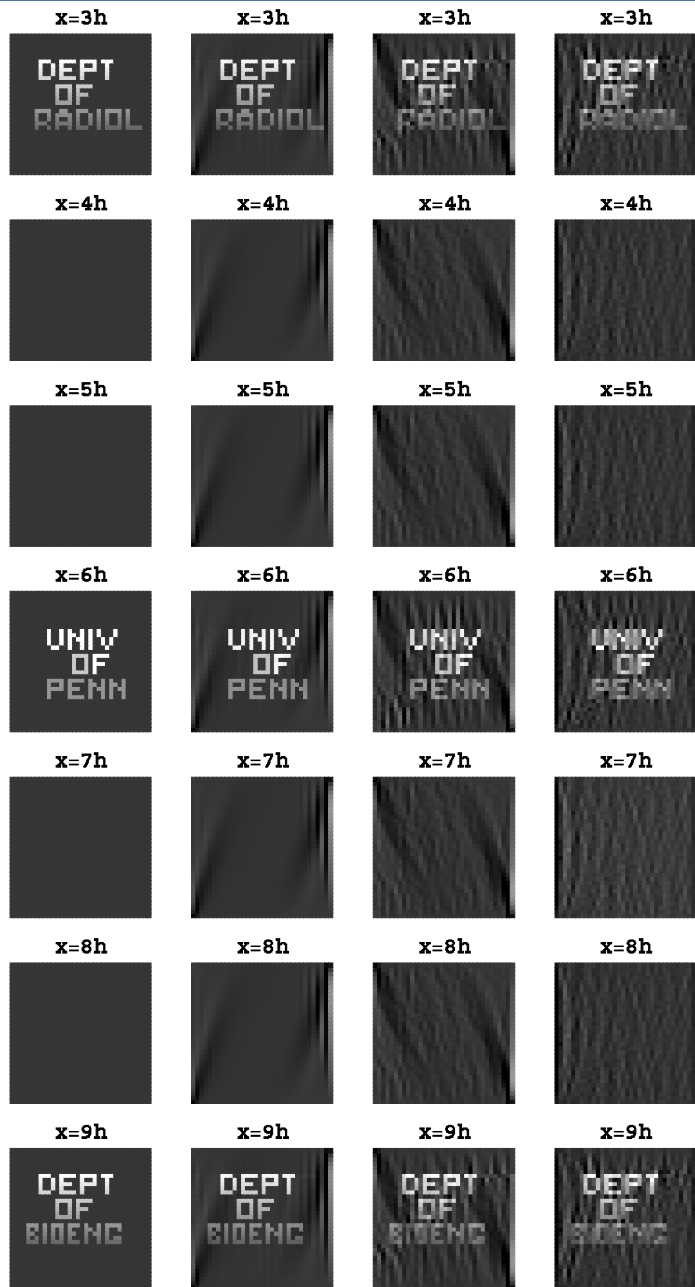
# OUTLINE

1. Motivation
2. Definition of the Star Transform
3. Physical Derivation of the Star Transform
4. Local Methods
5. Fourier Methods
6. Analysis of Stability
7. Numerical Examples

# 1. Motivation



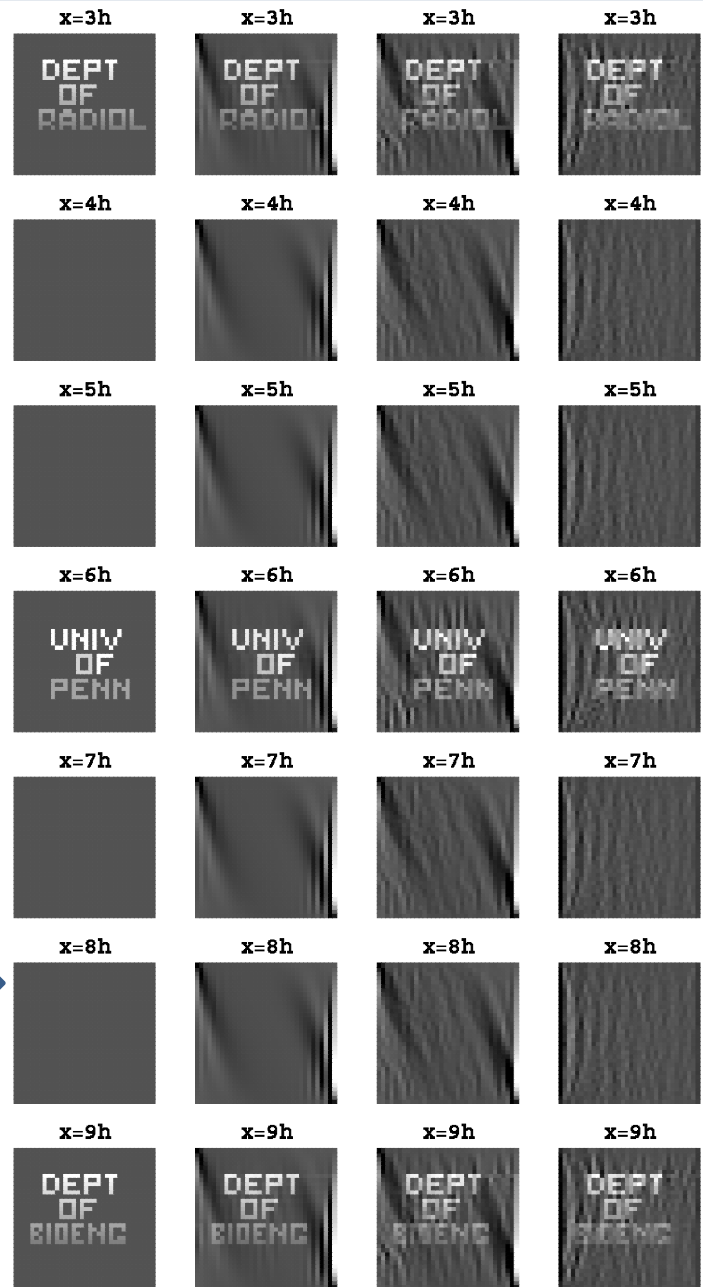
Model  $n = 0$   $n = 1\%$   $n = 3\%$  **Noise levels** Model  $n = 0$   $n = 1\%$   $n = 3\%$



$\mu_s h = 0.04$   
 $\mu_s L_z = 1.6$

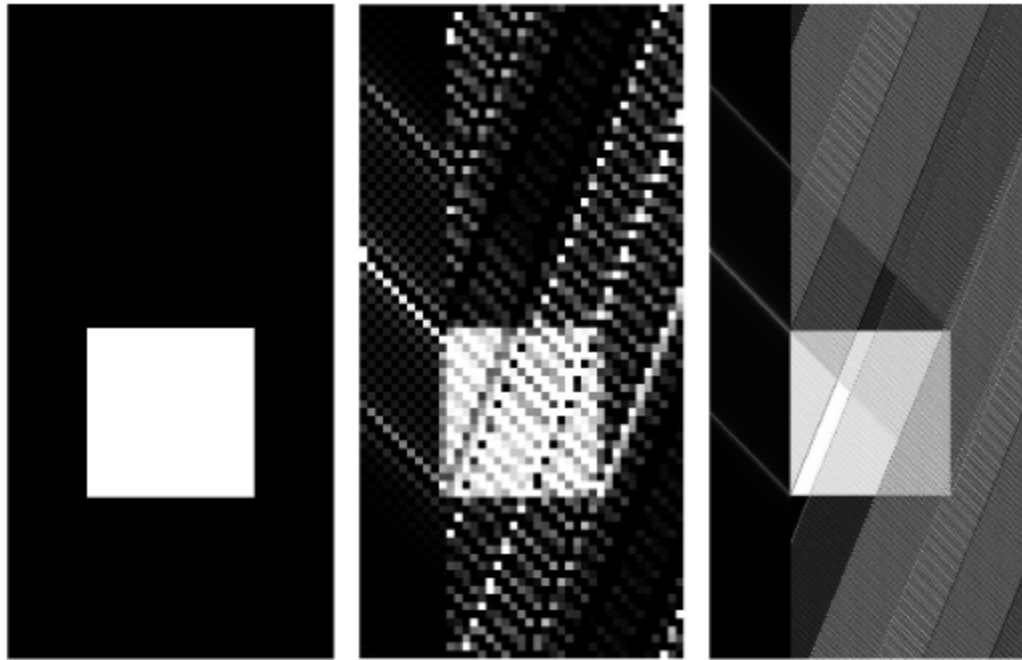
ABSORPTION  
 Background:  
 $\mu_a h = 0.01$   
 Targets:  
 $0.06 < \mu_a h < 0.2$

$\mu_s h = 0.08$   
 $\mu_s L_z = 3.2$



# Analytical reconstruction formula for one scattering angle (inverse-crime simulations)

$$\beta = \frac{\pi}{4}$$

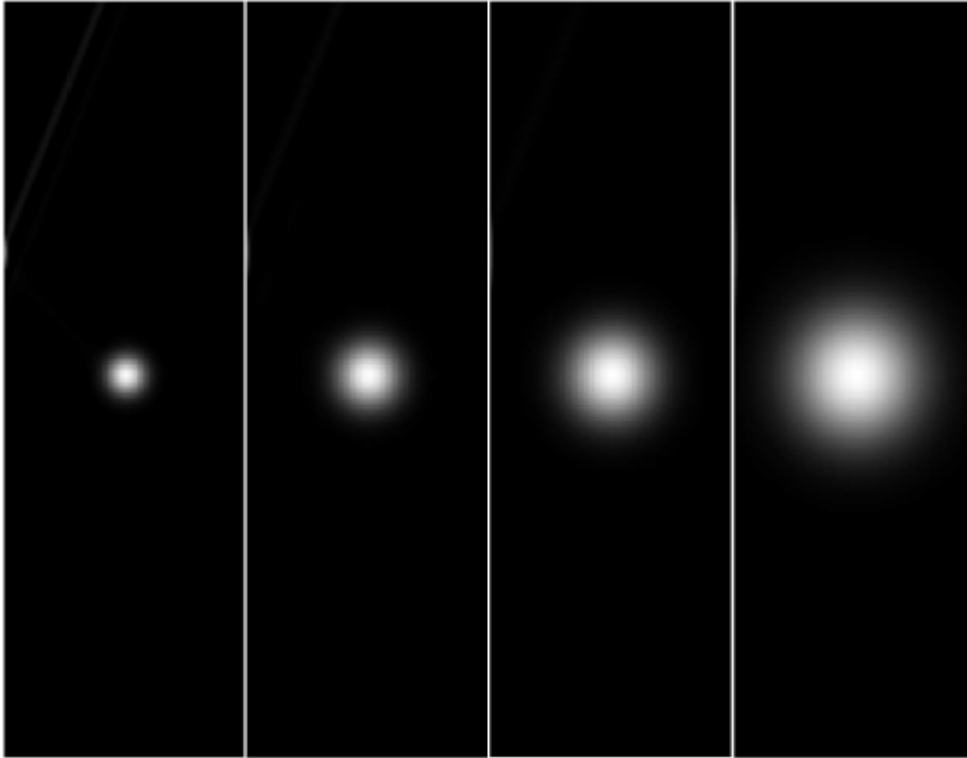


Model

$L/h=40$

$L/h=400$

# Reconstruction of a Gaussian



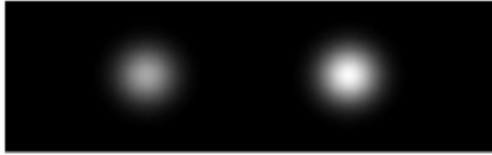
$$\beta = \frac{\pi}{4}$$
$$h = 0.3$$
$$\frac{L}{L} = \frac{40}{40}$$

$$\frac{w}{L} = \frac{3}{40}$$

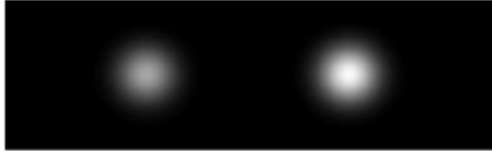
$$\frac{w}{L} = \frac{5}{40}$$

$$\frac{w}{L} = \frac{7}{40}$$

$$\frac{w}{L} = \frac{10}{40}$$

$\mu_t$  $\mu_s$  $\mu_a$ 

Model,  
 $w=7L/40$



Rec.



Model,  
 $w=3L/40$



Rec.



Model,  
 $w=1L/40$



Rec.

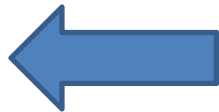
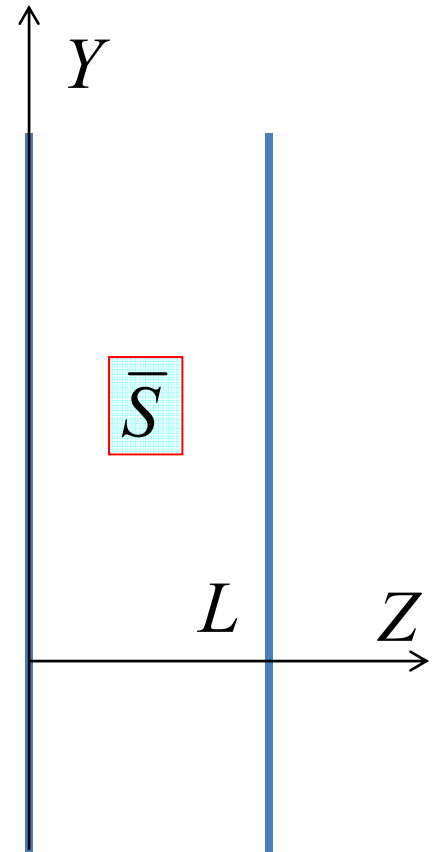
$$\frac{\bar{\mu}_s}{\bar{\mu}_a} = \frac{4}{3} ; \delta\mu_{a,s} = \bar{\mu}_{a,s} \exp\left[-\frac{(\mathbf{r} - \mathbf{r}_{a,s})^2}{w^2}\right] ; \frac{h}{L} = \frac{0.3}{40}$$

## 2. Definition of the Star Transform

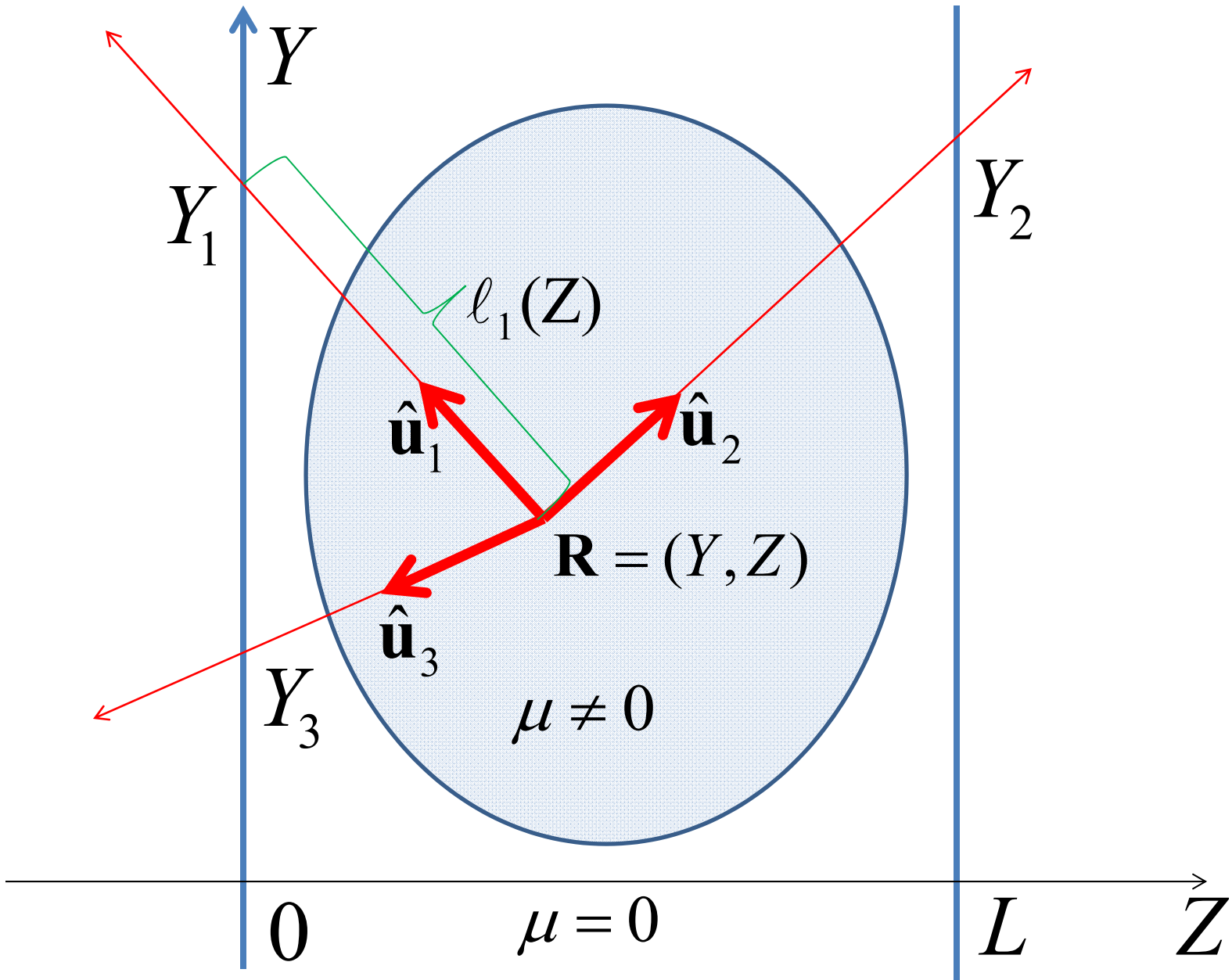
$$\Phi(\mathbf{R}) = \sum_{k=1}^K s_k I_k(\mathbf{R}) ,$$

$$\mathbf{R} \equiv (Y, Z) \in \bar{S} = \{0 \leq Z \leq L\} ,$$

$$I_k(\mathbf{R}) = \int_0^{\ell_k(Z)} \mu(\mathbf{R} + \hat{\mathbf{u}}_k \ell) d\ell$$







# 3. Physical Derivation of Star Transform

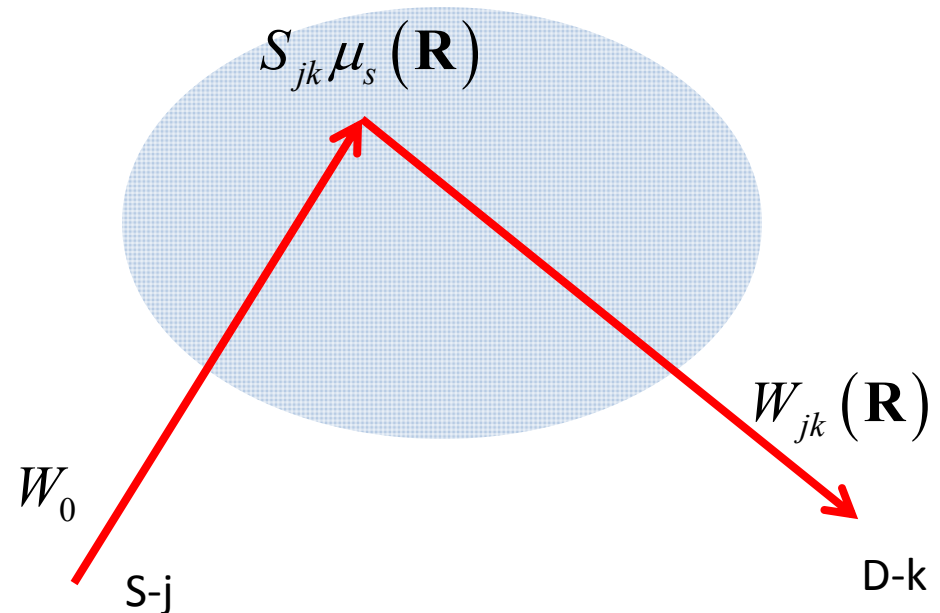
$$W_{jk}(\mathbf{R}) = W_0 S_{jk} \mu_s(\mathbf{R}) \exp[-I_j(\mathbf{R}) - I_k(\mathbf{R})]$$

$$\phi_{jk}(\mathbf{R}) = -\ln \left[ \frac{W_{jk}(\mathbf{R})}{W_0 S_{jk} \bar{\mu}_s} \right]$$

$$\phi_{jk}(\mathbf{R}) = I_j(\mathbf{R}) + I_k(\mathbf{R}) + \eta(\mathbf{R})$$

where

$$\eta(\mathbf{R}) = -\ln \left( \frac{\mu_s(\mathbf{R})}{\bar{\mu}_s} \right)$$



$$\phi_{jk}(\mathbf{R}) = \left[ I_j(\mathbf{R}) + I_k(\mathbf{R}) + \eta(\mathbf{R}) \right] (1 - \delta_{jk}) \quad (1)$$

$j, k = 1, 2, \dots, K$

$I_j(\mathbf{R})$  depend on total attenuation  $\mu(\mathbf{R})$

$\eta(\mathbf{R})$  depends on the scattering coefficient  $\mu_s(\mathbf{R})$

Strategy:

a) Exclude  $\eta(\mathbf{R})$  from the equations by considering linear combinations of  $\phi_{jk}(\mathbf{R})$ :

$$\Phi(\mathbf{R}) = \frac{1}{2} \sum_{j,k=1}^K c_{jk} \phi_{jk}(\mathbf{R}) = \sum_{k=1}^K s_k I_k(\mathbf{R})$$

b) Solve for total attenuation.

c) Using the above result, compute the ray integrals  $I_k(\mathbf{R})$ .

d) Use any of the equations in (1) to compute  $\eta(\mathbf{R})$ .

$$\Phi(\mathbf{R}) = \frac{1}{2} \sum_{j,k=1}^K c_{jk} \phi_{jk}(\mathbf{R})$$

Conditions on  $c_{jk}$

(i)  $c_{jk} = c_{kj}$

(ii)  $c_{kk} = 0$

(iii)  $\sum_{j,k=1}^K c_{jk} = 0$

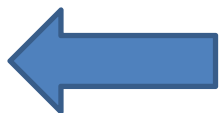
(iv)  $\sum_{j=1}^K c_{jk} = s_k \neq 0$

$$\Phi(\mathbf{R}) = \sum_{k=1}^K s_k I_k(\mathbf{R})$$

0	1	1		2
1	0	-2		-1
1	-2	0		-1
<hr/>				0
2	-1	-1		0

0	1	1	-1		1
1	0	-1	-1		-1
1	-1	0	1		1
-1	-1	1	0		-1
<hr/>					0
1	-1	1	-1		0

This last conditions is not critical.  
It excludes the possibility of star transforms  
in which a ray integral has zero "weight"



## 4. Local Methods

$$-(\hat{\mathbf{u}}_k \cdot \nabla) I_k(\mathbf{R}) = \mu(\mathbf{R})$$

Unfortunately, we can not make measurements of ray integrals  $I_k(\mathbf{R})$  directly. However, we can formulate the star transform so that the coefficients  $\mathbf{s}_k$  and  $\Phi$  are vectors. Then it is possible to invert the star transform by the local formula

$$\mu(\mathbf{R}) = \nabla \cdot \Phi(\mathbf{R})$$

[Katsevich and Krylov, Inverse Problems **29**, 075008 (2013)]

What if we allow the coefficients  $\mathbf{c}_{jk}$  to be vectors?

Moreover, let

$$\sum_{j=1}^K \mathbf{c}_{jk} = \mathbf{s}_k = \sigma_k \hat{\mathbf{u}}_k \quad \text{and} \quad \sum_{k=1}^K \mathbf{s}_k = \sum_{k=1}^K \sigma_k \hat{\mathbf{u}}_k = \mathbf{0}$$

Then define

$$\Phi(\mathbf{R}) \equiv \frac{1}{2} \sum_{j,k=1}^K \mathbf{c}_{jk} \phi_{jk}(\mathbf{R})$$

and

$$\Phi(\mathbf{R}) = \sum_{k=1}^K \sigma_k \hat{\mathbf{u}}_k I_k(\mathbf{R})$$

$$-\nabla \cdot \Phi(\mathbf{R}) = \left( \sum_{k=1}^K \sigma_k \right) \mu(\mathbf{R})$$

$$\mu(\mathbf{R}) = -\frac{1}{\sum_{k=1}^K \sigma_k} \nabla \cdot \Phi(\mathbf{R})$$

$\mathbf{c}_{jk}$  matrix for  $K = 3$

0	$\sigma_1 \hat{\mathbf{u}}_1 + \sigma_2 \hat{\mathbf{u}}_2$	$\sigma_1 \hat{\mathbf{u}}_1 + \sigma_3 \hat{\mathbf{u}}_3$	$\sigma_1 \hat{\mathbf{u}}_1$
$\sigma_1 \hat{\mathbf{u}}_1 + \sigma_2 \hat{\mathbf{u}}_2$	0	$\sigma_2 \hat{\mathbf{u}}_2 + \sigma_3 \hat{\mathbf{u}}_3$	$\sigma_2 \hat{\mathbf{u}}_2$
$\sigma_1 \hat{\mathbf{u}}_1 + \sigma_3 \hat{\mathbf{u}}_3$	$\sigma_2 \hat{\mathbf{u}}_2 + \sigma_3 \hat{\mathbf{u}}_3$	0	$\sigma_3 \hat{\mathbf{u}}_3$
$\sigma_1 \hat{\mathbf{u}}_1$	$\sigma_2 \hat{\mathbf{u}}_2$	$\sigma_3 \hat{\mathbf{u}}_3$	0

Here  $\sigma_k$  are chosen so that  $\sigma_1 \hat{\mathbf{u}}_1 + \sigma_2 \hat{\mathbf{u}}_2 + \sigma_3 \hat{\mathbf{u}}_3 = 0$

Reconstruction formula:

$$\mu = -\frac{1}{\sigma_1 + \sigma_2 + \sigma_3} \nabla \cdot \left[ \sigma_1 \hat{\mathbf{u}}_1 (\phi_{12} + \phi_{13}) + \sigma_2 \hat{\mathbf{u}}_2 (\phi_{21} + \phi_{23}) + \sigma_3 \hat{\mathbf{u}}_3 (\phi_{31} + \phi_{32}) \right]$$

Symmetric case  $\hat{\mathbf{u}}_1 + \hat{\mathbf{u}}_2 + \hat{\mathbf{u}}_3 = 0$ :

$$\mu = -\frac{1}{3} \nabla \cdot \left[ \hat{\mathbf{u}}_1 (\phi_{32} - \phi_{21}) + \hat{\mathbf{u}}_2 (\phi_{31} - \phi_{12}) \right]$$

$c_{jk}$  matrix for  $K = 4$  (method of Katsevich and Krylov)

$$\sigma_1 = 0, \sigma_2 = 1, \sigma_3 = -a, \sigma_4 = -b$$

$$-a - b = 0$$

$$\hat{\mathbf{u}}_2 - a\hat{\mathbf{u}}_3 - b\hat{\mathbf{u}}_4 = 0$$

$$\begin{array}{cccc|c} 0 & \hat{\mathbf{u}}_2 & -a\hat{\mathbf{u}}_3 & -b\hat{\mathbf{u}}_4 & 0 \\ \hat{\mathbf{u}}_2 & 0 & 0 & 0 & \hat{\mathbf{u}}_2 \\ -a\hat{\mathbf{u}}_3 & 0 & 0 & 0 & -a\hat{\mathbf{u}}_3 \\ -b\hat{\mathbf{u}}_4 & 0 & 0 & 0 & -b\hat{\mathbf{u}}_4 \\ \hline 0 & \hat{\mathbf{u}}_2 & -a\hat{\mathbf{u}}_3 & -b\hat{\mathbf{u}}_4 & 0 \end{array}$$

Reconstruction formula:

$$\mu = \frac{1}{a+b-1} \nabla \cdot [\hat{\mathbf{u}}_2 \phi_{12} - a\hat{\mathbf{u}}_3 \phi_{13} - b\hat{\mathbf{u}}_4 \phi_{14}]$$



# 5. Fourier Methods

$$\Phi(\mathbf{R}) = \sum_{k=1}^K s_k I_k(\mathbf{R}) \quad , \quad I_k(\mathbf{R}) = \int_0^{\ell_k(Z)} \mu(\mathbf{R} + \hat{\mathbf{u}}_k \ell) d\ell$$

$$\mu(y, z) = \int_{-\infty}^{\infty} \frac{dq}{2\pi} e^{iqy} \tilde{\mu}(q, z) = \int_{-\infty}^{\infty} \frac{dq}{2\pi} e^{iqy} \frac{1}{L} \sum_{n=-\infty}^{\infty} \mu_n(q) e^{i\kappa_n z} \quad ,$$

$$\mu_n(q) = \int_{-\infty}^{\infty} dy e^{-iqy} \int_0^L dz e^{-i\kappa_n z} \mu(y, z) \quad ,$$

$$\kappa_n = \frac{2\pi n}{L} \quad .$$

Fix  $q$ . For each  $q$ , we obtain a set of linear equations  
(different  $q$ 's are not mixed)

$$d_n \mu_n + \sum_{k=1}^K s_k \alpha_k \frac{1}{\beta_k + \kappa_n} \sum_{m=-\infty}^{\infty} \frac{\mu_m}{\beta_k + \kappa_m} = \Phi_n$$

$$\alpha_k = \frac{e^{i\beta_k \xi_k} (e^{-i\beta_k L} - 1)}{L u_{kz}}, \quad \xi_k = 0, L$$

$$\beta_k = q \frac{u_{ky}}{u_{kz}}$$

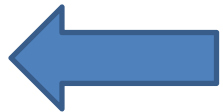
$$d_n = \sum_{k=1}^K \frac{i s_k}{u_{kz} (\beta_k + \kappa_n)} = \sum_{k=1}^K \frac{i s_k}{\hat{\mathbf{u}}_k \cdot (q, \kappa_n)}$$

In matrix notations:

$$D|\mu\rangle + \sum_{k=1}^K s_k \alpha_k |\psi_k\rangle \langle \psi_k | \mu \rangle = |\Phi\rangle$$

$D$  is diagonal

$|\psi_k\rangle$  are known vectors



# 6. Analysis of Stability

- a)  $qL \ll 1$  : Can be easily analyzed
- b)  $qL \gg 1$  : Main concern. Not so simple but can also be analyzed
- c)  $qL \sim 1$  : No analytical condition obtained; empirically, we have found that, if all rays cross the same boundary (reflection geometry), there can exist an instability in reconstructions. However, corresponding artifacts are localized near the the boundaries and are not of major concern.

## a) Case $q = 0$

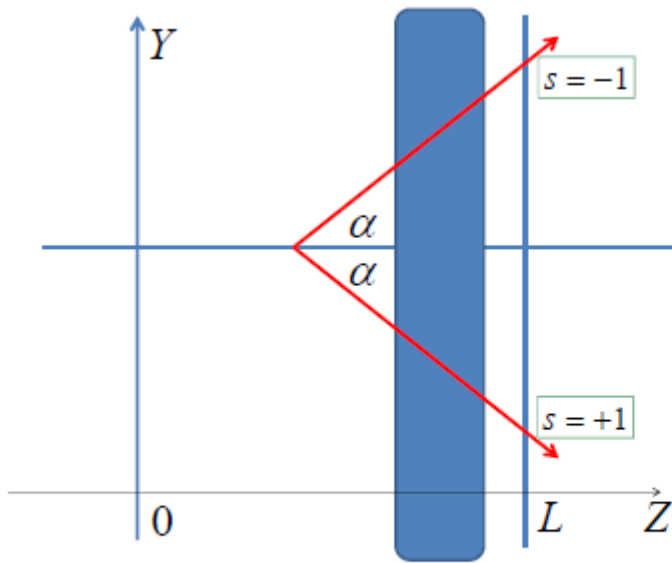
$$\mu_0 = \frac{2}{L\Sigma_0} \sum_{m=-\infty}^{\infty} \Phi_m = \frac{2\tilde{\Delta}(0)}{\Sigma_0}, \quad n = 0,$$

$$\mu_n = \mu_0 - i \frac{\kappa_n \Phi_n}{\Sigma_1}, \quad n \neq 0.$$

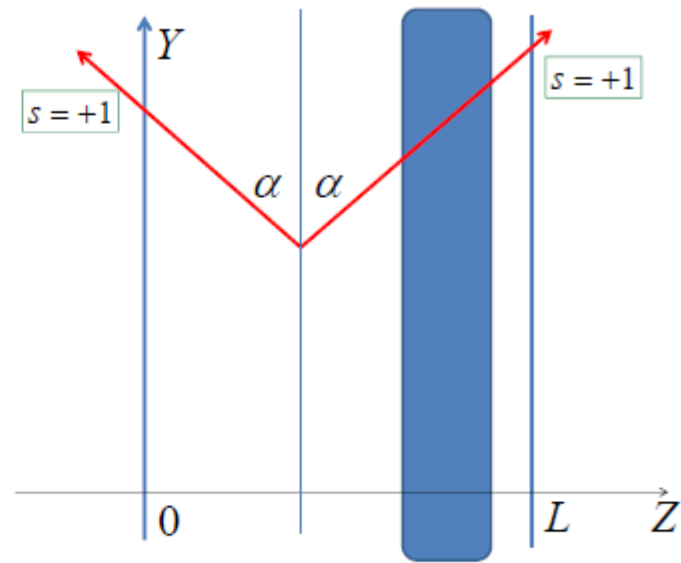
$$\Sigma_0 = \sum_{k=1}^K \frac{s_k}{|u_{kz}|}, \quad \Sigma_1 = \sum_{k=1}^K \frac{s_k}{u_{kz}}, \quad \Sigma_2 = - \sum_{k=1}^K \frac{s_k u_{ky}}{u_{kz}^2}.$$

$$d_n = i \left( \Sigma_1 \frac{1}{\kappa_n} + \Sigma_2 \frac{q}{\kappa_n^2} + \dots \right), \quad |n| \rightarrow \infty,$$

# Bad imaging geometries ( $\Sigma_1=0$ )



(a)



(b)

b) Case  $q \rightarrow \infty$

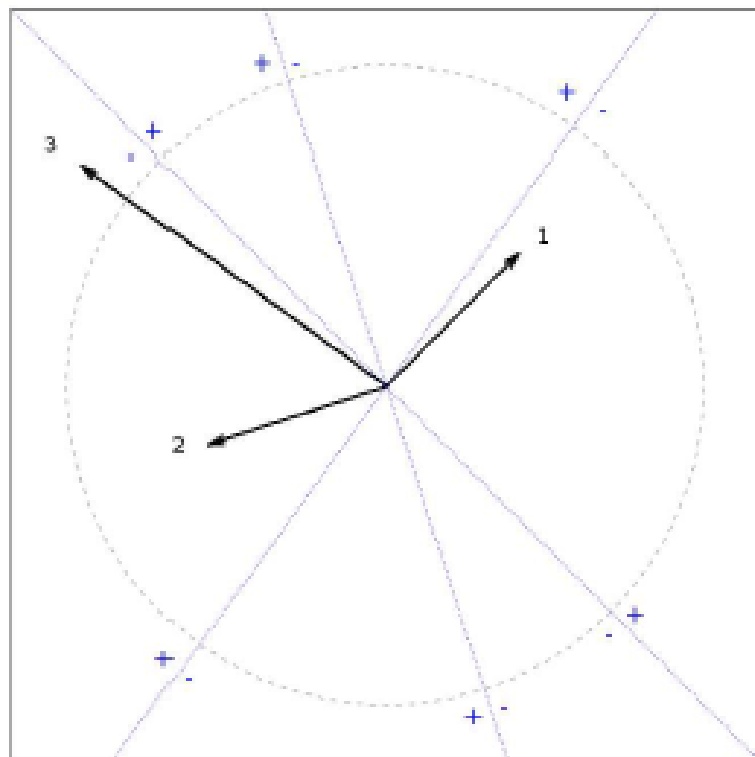
The diagonal matrix  $D$  somonates the separable terms

Find the condition under which all elements  $d_n$  of  $D$  are simultaneously non-zero

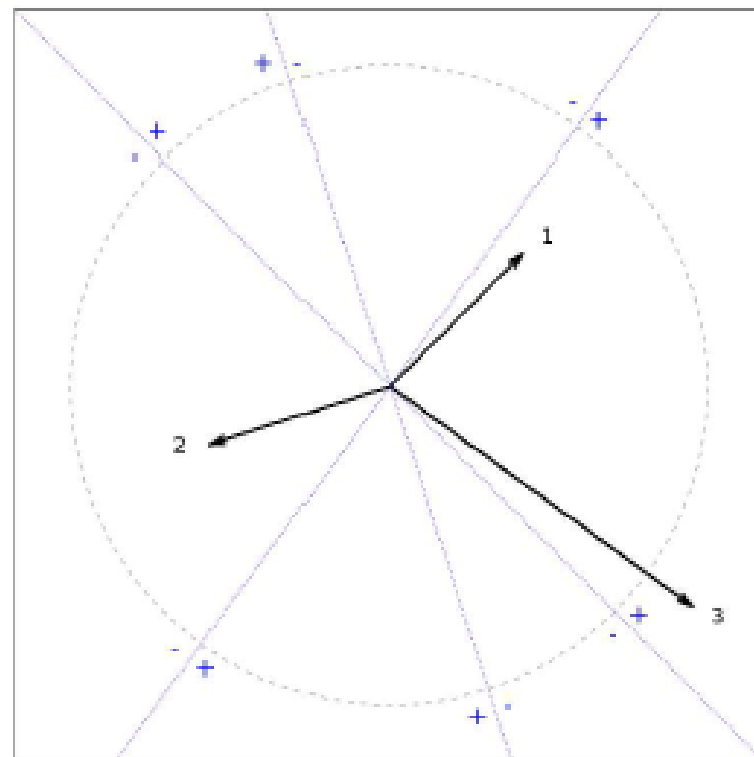
$$d_n(q) = \sum_{k=1}^K \frac{i s_k}{\hat{\mathbf{u}}_k \cdot (q, \kappa_n)} = \frac{i}{|(q, \kappa_n)|} \sum_{k=1}^K \frac{s_k}{\hat{\mathbf{u}}_k \cdot \hat{\mathbf{v}}}$$

$$f(\theta) = \sum_{k=1}^K \frac{s_k}{\cos(\theta - \theta_k)}$$

$f(\theta)$  has zeros  $\Leftrightarrow$  at least one of the elements  $d_n(q)$  turns to zero for some  $q$  and  $n$ .  
There exists a high-frequency instability

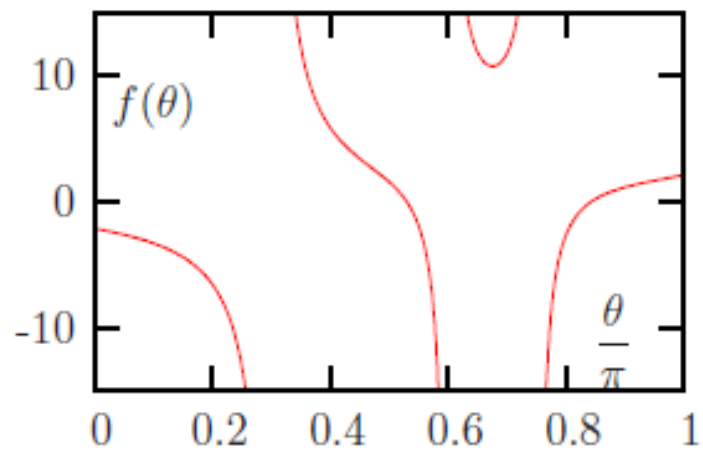


(a)

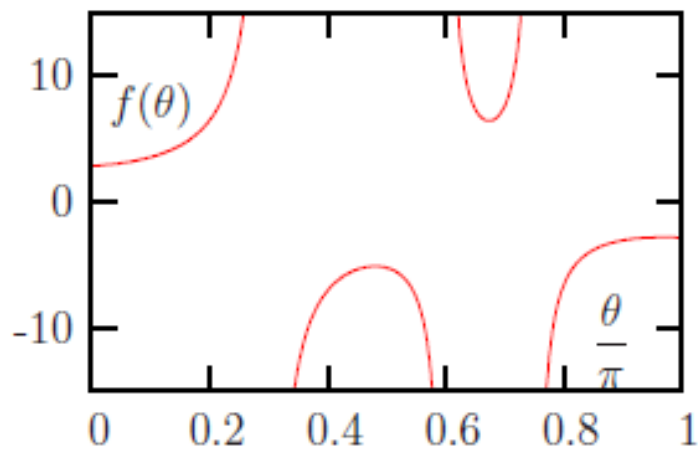


(b)

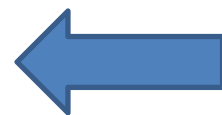
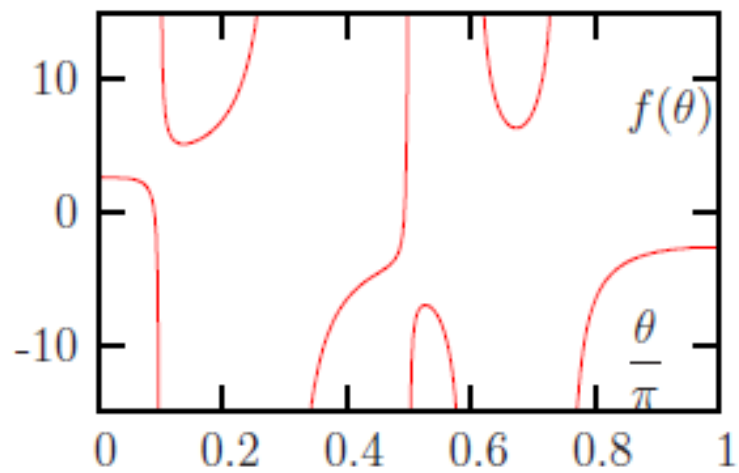




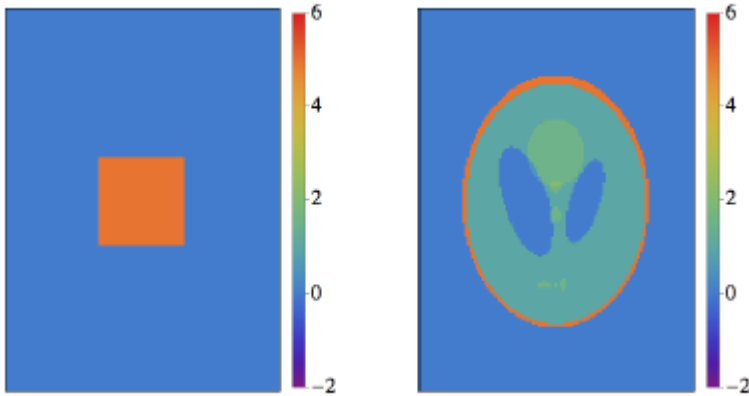
(a)



(b)



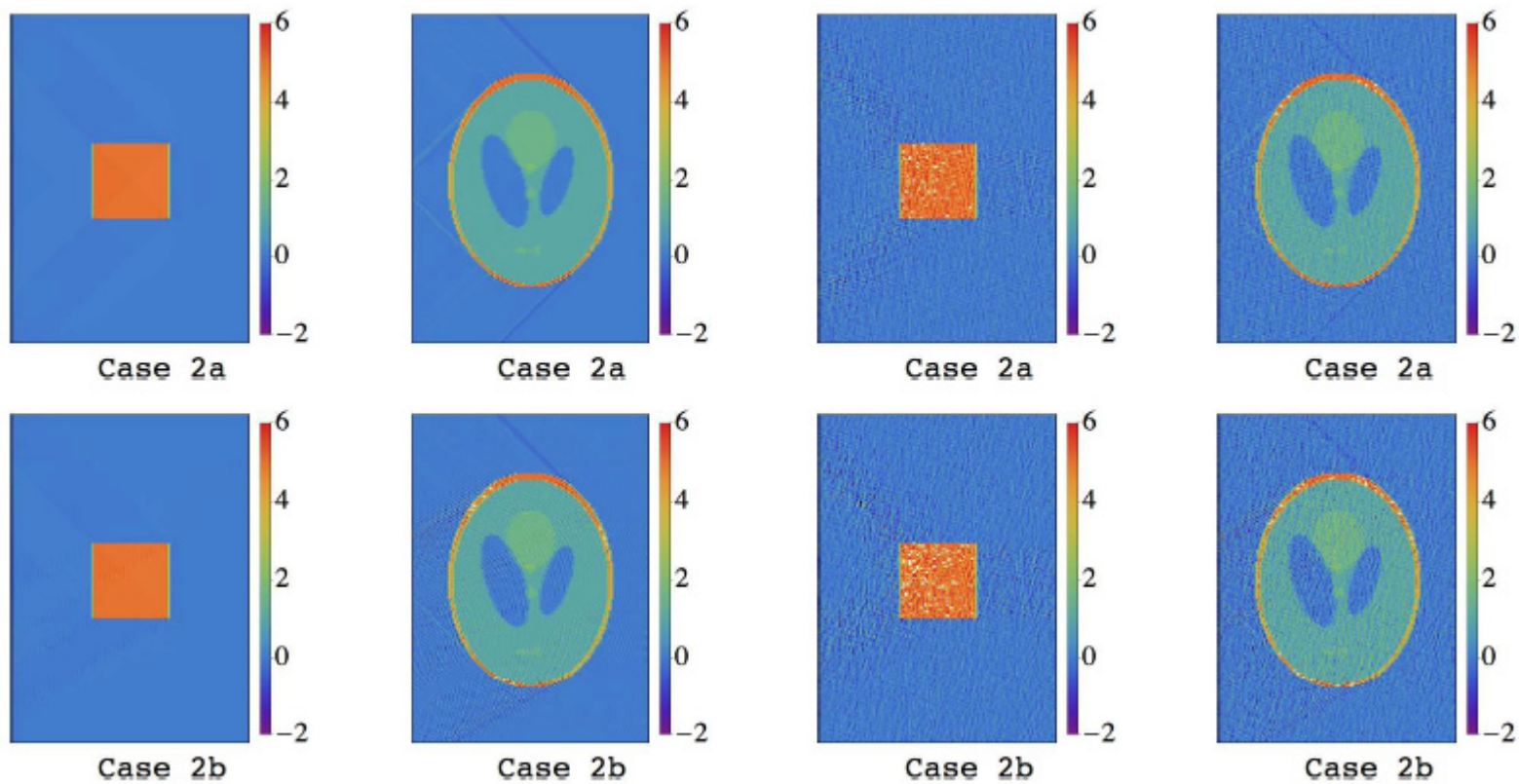
# 7. Numerical Examples



Phantoms used

Case	2a	2b	3a	3b
K	3		3	
$s_1$	1		1	
$s_2$	1		1	
$s_3$	1		-2	
$\theta_1/\pi$	1		0.25	
$\theta_2/\pi$	0.25		1.1	
$\theta_3/\pi$	-0.25	-1/6	-0.2	0.8
NZ	0	0	2	0
$\Sigma_0$	3.83	3.57	-0.01	-0.01
$\Sigma_1$	1.83	1.57	-2.11	2.83

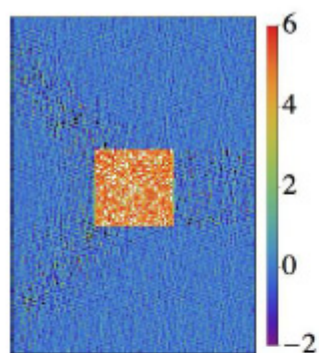
CASE 2;  $K=3$ ,  $s_1=s_2=s_3=1$



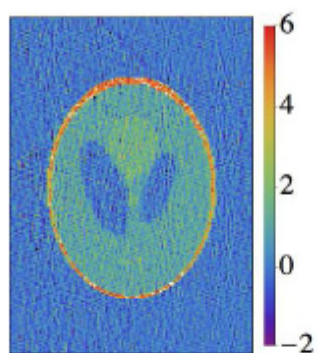
(a) No noise,  $\lambda = 0$

(b)  $\mathcal{N} = 4 \times 10^4$ ,  $\lambda = 0$

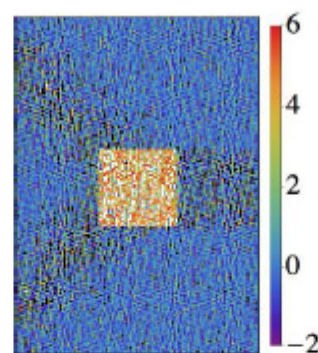
CASE 2;  $K=3$ ,  $s_1=s_2=s_3=1$



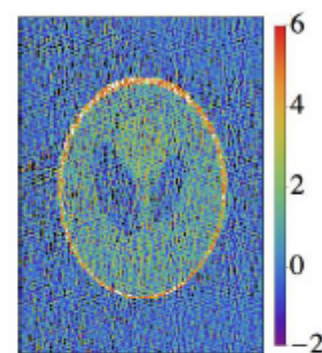
Case 2a



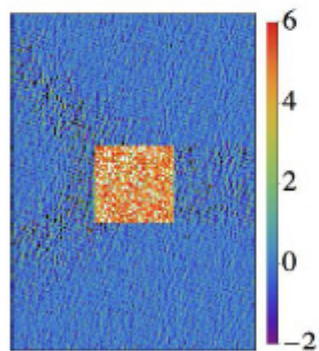
Case 2a



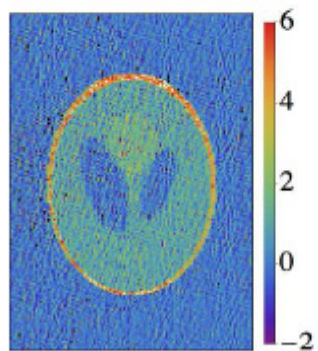
Case 2a



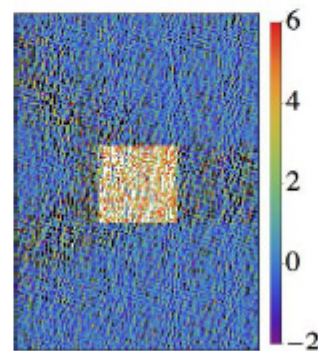
Case 2a



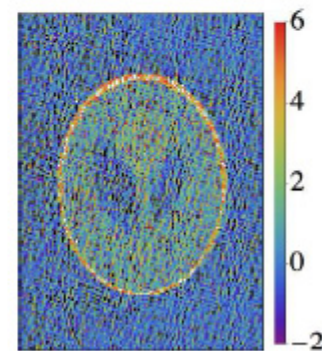
Case 2b



Case 2b



Case 2b

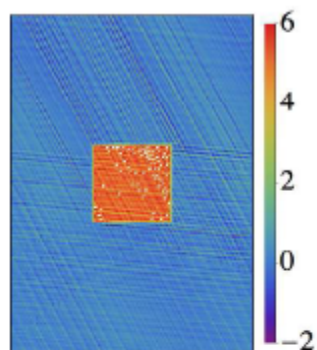


Case 2b

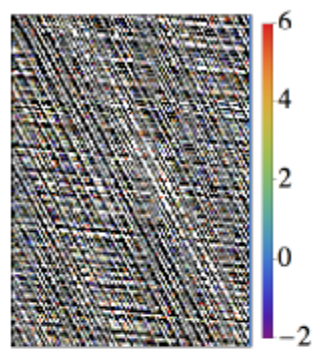
(c)  $\mathcal{N} = 10^4$ ,  $\lambda = 0$

(d)  $\mathcal{N} = 2.5 \times 10^3$ ,  $\lambda = 0$

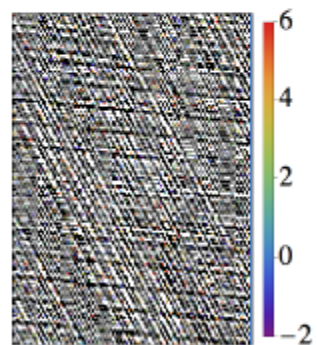
CASE 3;  $K=3$ ,  $s_1=s_2=1$ ,  $s_3=-2$



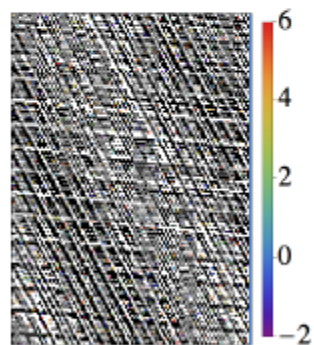
No Noise



$N=4 \times 10^4$

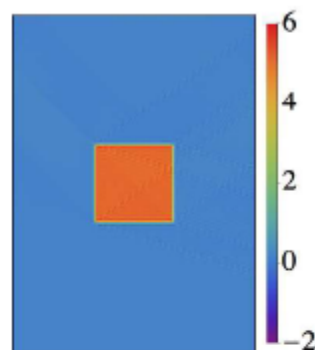


$N=10^4$

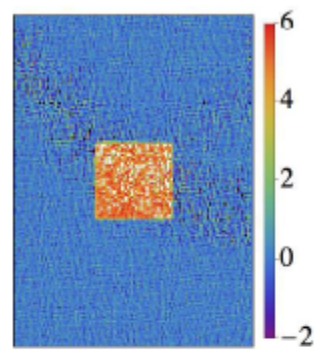


$N=2.5 \times 10^3$

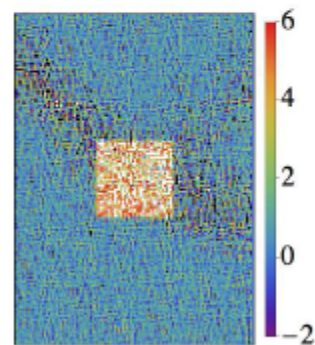
(a) Case 3a



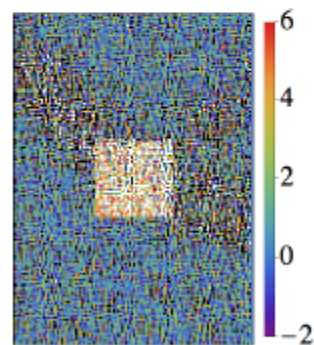
No Noise



$N=4 \times 10^4$



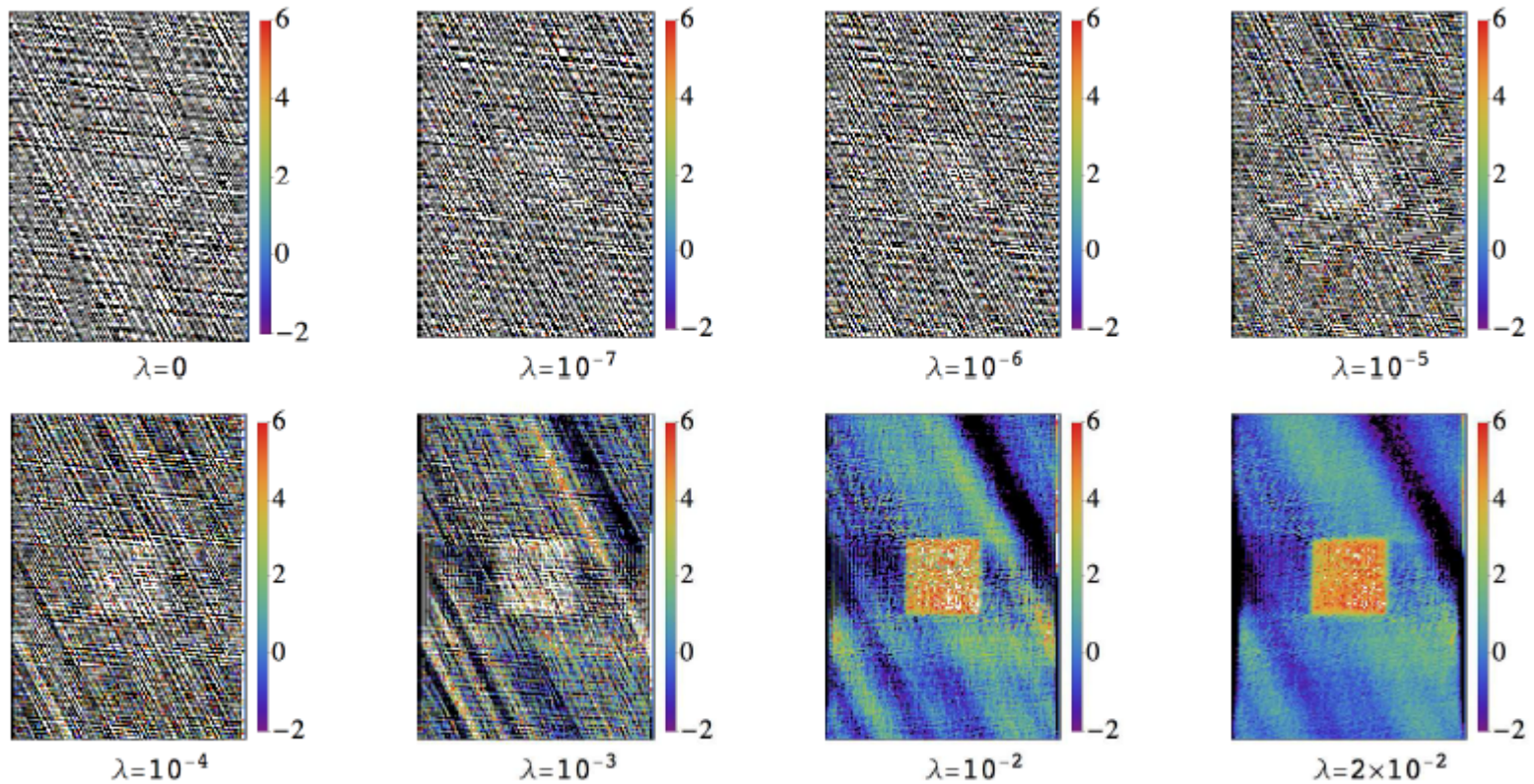
$N=10^4$



$N=2.5 \times 10^3$

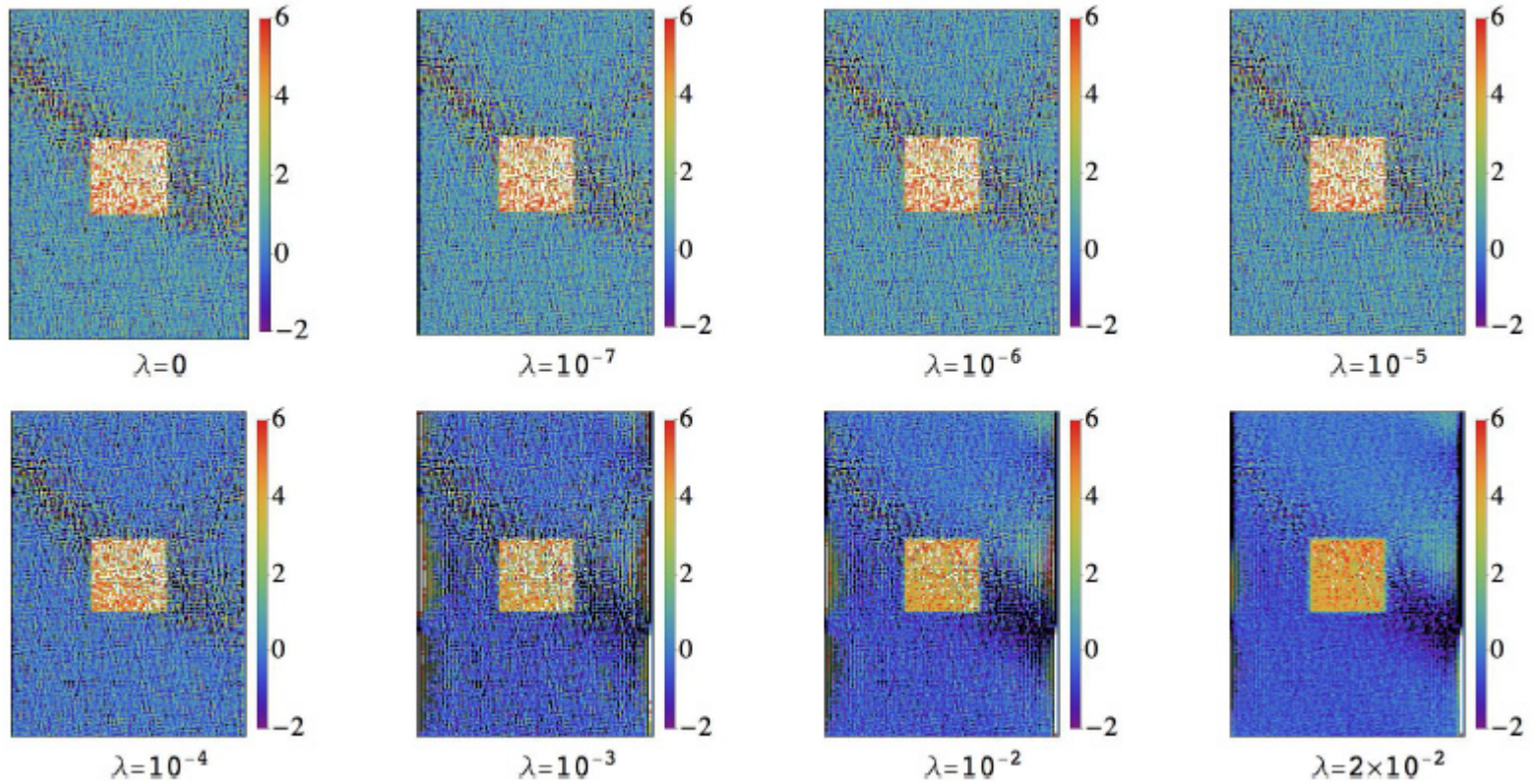
(b) Case 3b

CASE 3a;  $K=3$ ,  $s_1=s_2=1$ ,  $s_3=-2$ ,  $NZ=2$ ; effects of regularization for  $N=1e4$



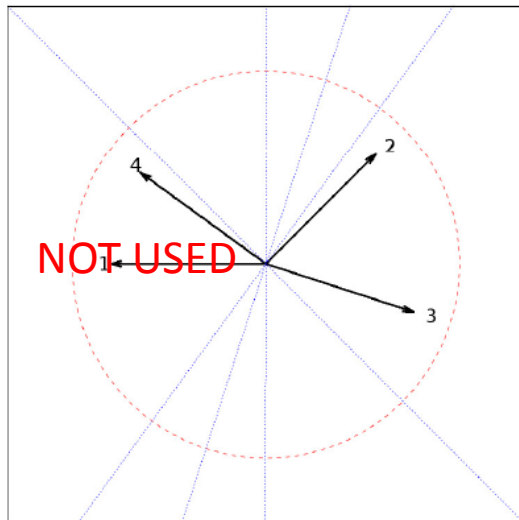
(a) Case 3a

CASE 3b;  $K=3$ ,  $s_1=s_2=1$ ,  $s_3=-2$ ,  $NZ=0$ ; effects of regularization for  $N=1e4$

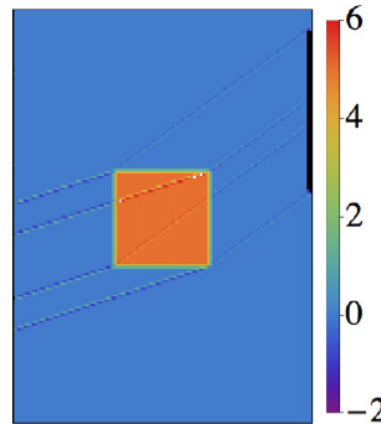


(b) Case 3b

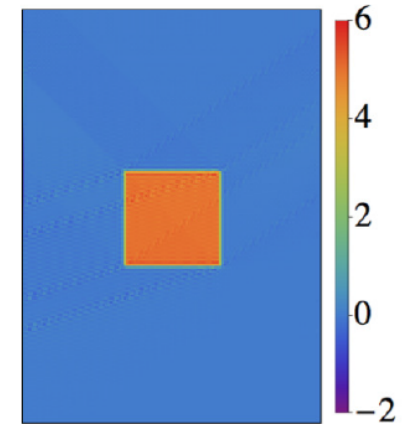
# Comparison of local and Fourier methods (K=3 coefficient matrix)



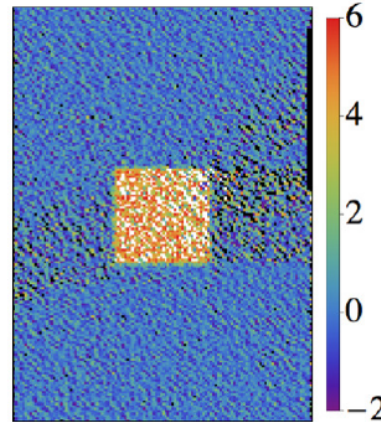
	$\alpha_k/\pi$	$C_k$	$S_k$	$\sigma_k$
1	1.0	NOT USED		
2	0.25	1	1	-0.3468
3	-0.1	-2	-2	1.1085
4	0.8	1	1	1.0000



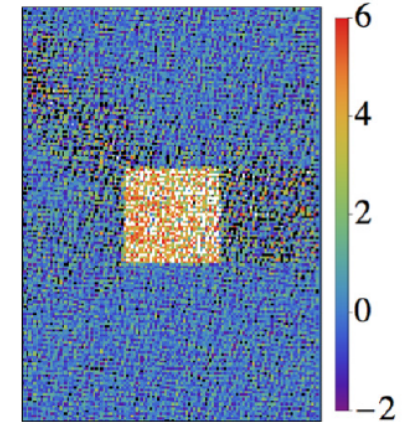
Local, without noise



Fourier, without noise



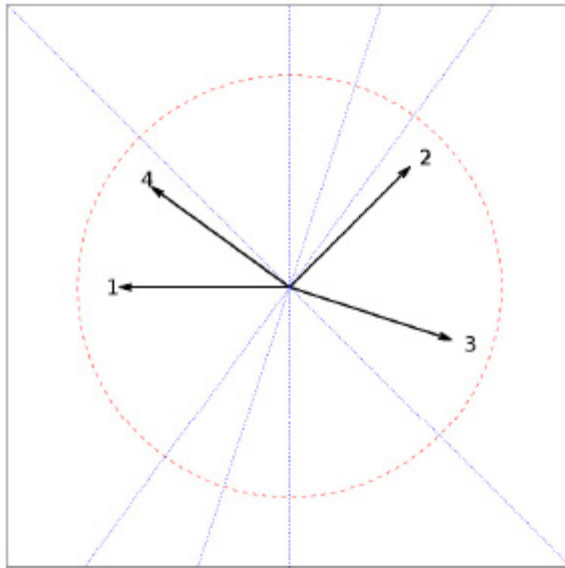
Local, with noise



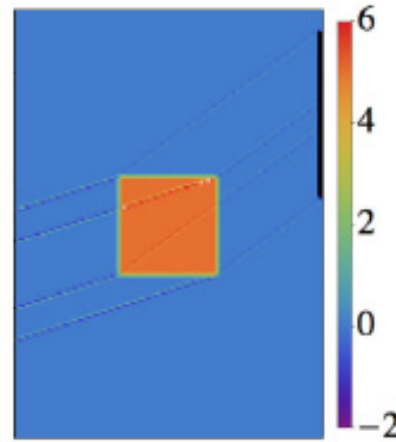
Fourier, with noise



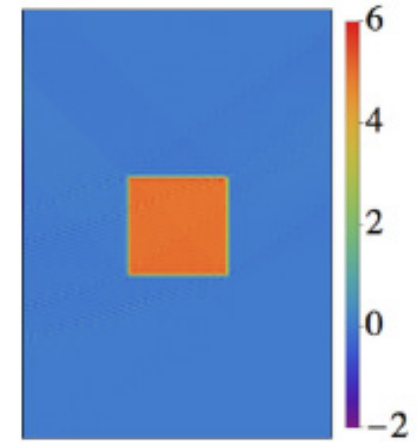
Comparison of local and Fourier methods (K=4, geometry of Katsevich and Krylov, first ray canceled in the star transform)



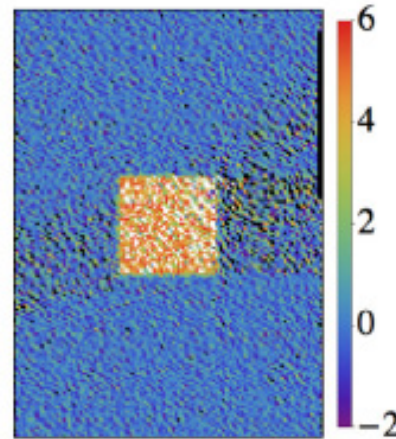
	$\alpha_k/\pi$	$c_k$	$s_k$	$\sigma_k$
1	1.0			
2	0.25	1	1	-0.3468
3	-0.1	-2	-2	1.1085
4	0.8	1	1	1.0000



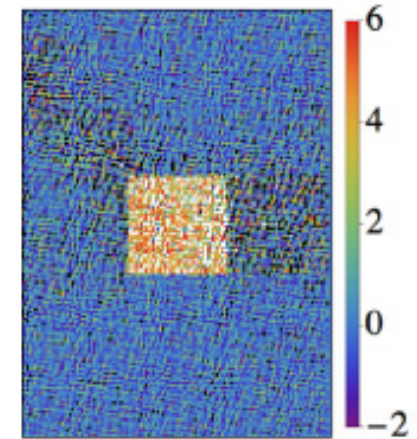
Local, without noise



Fourier, without noise



Local, with noise



Fourier, with noise